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Photonics, Particle Accelerators and beyond the Standard Model



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Figure 1 Schéma de la distribution du champ électrique dans ani cavité résonnante















Photonic

Metamaterial

a << λ

 $a \sim \lambda$









ERSF 75KW Solid State RF



Hughes TWT - RF Tube 100kW



ancaster University





FIG. 1: The traveling wave structure **considered** here, consisting of a folded waveguide with a metamaterial insert, the electron beam passes through the middle of the structure $\omega = \sqrt{c^2 \left(\gamma_n - \left(\frac{\pi(2n+1)}{p} + \beta_{mm}(f)\frac{\Delta h}{p}\right)\right)} \alpha + \omega_c^2$

metamaterial, to define a unique beam-wave interaction, triggering 0 hovel gain-frequency phenomena + $(2n+1)\frac{\pi}{p}$

$$\alpha = \left(\frac{\mathbf{H} p_{\mathbf{METH}}}{p + h - \Delta h}\right)^2 \mathbf{DOLOGY}$$

For effective energy transfer between beam and EM wave the phase velocity (determined by the dispersion) of the wave must approximately match the velocity of the electron beam. In the conventional FWTWT this is achieved (ia) the periodicity of the (fo)) ad (vaveguide $\frac{2}{c}$ [19] to slow down the wave, generating Spatial Harmonics Wave Components (SHWC) parallel to the beam. The SHWC interact with the beam resulting in energy transfer. By a superposition of the spatial harmonics $(\vec{E}_m(x, y))$ the field parallel to the beam can be expressed by Floquets theorem [20] as;

$$\vec{E}(z) = \sum_{m=-\infty}^{\infty} \vec{E}_m(x, y) e^{-i\beta_m z}, \beta_m = \beta + \frac{2m\pi}{p} \qquad (1)$$





Lorentz's Force Equation



Time changing in $m_0 c^2 \gamma$ (DC and AC beam energy) is related by the E.v dot product in this equation. The DC beam energy γdc is given by (1+Vdc/511); Vdc is the DC beam accelerating potential. While the AC beam energy exchange (stimulated emission) is calculated through the Madey's theory.

1st order perturbation >> Spontaneous emission

 $\left\langle \Delta \gamma_2 \right\rangle = \frac{1}{2} \frac{d}{d} \left\langle \Delta \gamma_1^2 \right\rangle$ 2nd order perturbation >> Stimulated emission







Complementary Split Ring Resonator







Transmission through CSRR interrupting waveguide.







- 1.6 KeV/m Acc gradient
- About 1/5 of the gradient to a comparable pill box resonator

Dispersion relation extracted, black dots, with the light line shown in green.





























Xu and Seviour 2012 New J. Phys. 14 013014



Yoon Kang, Ali Nassiri, IPAC 1998





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$$\begin{split} \vec{H}_{ij} &= \langle \vec{E}_0(\vec{r} - \vec{R}_i) | \hat{O} | \vec{E}_0(\vec{r} - \vec{R}_j) \rangle \\ &= \begin{cases} (\omega_0/c)^2 = \alpha \quad (i = j), \text{the same defect} \\ (\omega_0/c)^2 \beta_1 \quad (i \neq j), \text{the first-neighbour defect} \\ (\omega_0/c)^2 \beta_2 \quad (i \neq j), \text{the second-neighbour defect} \\ (\omega_0/c)^2 \beta_3 \quad (i \neq j), \text{the third-neighbour defect} \end{cases} \begin{pmatrix} \alpha - \gamma \quad \tilde{\beta}_1 \quad \tilde{\beta}_2 \quad \tilde{\beta}_3 & \tilde{\beta}_2 \\ \tilde{\beta}_2 \quad \tilde{\beta}_1 \quad \alpha - \gamma \quad \tilde{\beta}_1 & \tilde{\beta}_2 & \tilde{\beta}_3 \\ \tilde{\beta}_3 \quad \tilde{\beta}_2 \quad \tilde{\beta}_1 \quad \alpha - \gamma \quad \tilde{\beta}_1 & \tilde{\beta}_2 \\ \tilde{\beta}_2 \quad \tilde{\beta}_3 \quad \tilde{\beta}_2 \quad \tilde{\beta}_1 \quad \alpha - \gamma \quad \tilde{\beta}_1 & \tilde{\beta}_2 \\ \tilde{\beta}_2 \quad \tilde{\beta}_3 \quad \tilde{\beta}_2 \quad \tilde{\beta}_1 \quad \alpha - \gamma \quad \tilde{\beta}_1 \\ \tilde{\beta}_1 \quad \tilde{\beta}_2 \quad \tilde{\beta}_3 \quad \tilde{\beta}_2 \quad \tilde{\beta}_1 \quad \alpha - \gamma \quad \tilde{\beta}_1 \\ \tilde{\beta}_1 \quad \tilde{\beta}_2 \quad \tilde{\beta}_3 \quad \tilde{\beta}_2 \quad \tilde{\beta}_1 \quad \alpha - \gamma \quad \tilde{\beta}_1 \end{cases} \end{split}$$

11-11

= 0

 $\alpha - \gamma$















Hidden

Sector Photon Searches





$$\nabla^2 E - \mu \varepsilon_r \varepsilon_0 \omega^2 \frac{\partial^2 E}{\partial t^2} = 0$$





$$\nabla^2 E - \mu \varepsilon \cdot \omega^2 \frac{\partial^2 E}{\partial t^2} = 0$$





$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} - \frac{\chi}{2} F^{\mu\nu} B_{\mu\nu} + \frac{m_{\gamma}^2}{2} B^{\mu} B_{\mu}$$









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